

# Study of Normalized Moments of Relativistic Charged Shower Particles Produced in $^{12}\text{C}$ -Nucleus Collisions at 4.5 A GeV

Praveen Prakash Shukla<sup>1</sup>, H. Khushnood<sup>2</sup>, M. Saleem Khan<sup>3</sup>

<sup>1,3</sup>Department of Applied Physics, MJPR, University Bareilly-243001

<sup>2</sup>University Polytechnic, Jamia Millia Islamia New Delhi-110025

**Abstract:** An attempt has been made to study the characteristics of charged shower particles produced in  $^{12}\text{C}$ -nucleus interactions at 4.5 A GeV. Results shows that the values of normalized moments of the multiplicity distribution of relativistic charged particles produced in relativistic energy in proton-nucleus and nucleus-nucleus reactions are almost constant within the statistical limit. Furthermore, it is reported that the leading particle multiplicity strongly depends on the mass of the projectile.

**Keywords:** Normalized moments, multiplicity distributions, charged shower particles, leading particle etc.

## 1. Introduction

The study of the characteristics of secondary charged particles produced in nucleus-nucleus reactions at relativistic energies has received considerable attention during recent years [1-6]. The reasons for investigating the production mechanism of secondary particles in heavy ion interaction might provide some valuable information about reaction mechanism in the nucleus-nucleus collisions. In relativistic heavy ion collisions, the secondary hadrons are formed instantaneously. There is a formation time between the collision and hadronization of the final state particles. They hadronize within the target nucleus and may re-interact with the surrounding target matter and produce cascade particles. Moreover, the studies on heavy ion interactions provide the means of discrimination between the various theoretical models put forward to explain the mechanism of hadronization of final state charged particles in nucleus-nucleus reactions at relativistic energies and the characteristics of secondary charged particles produced in such collisions may be obtained by analyzing the experimental data on the secondary charged particles.

## 2. Experimental Technique

In the present investigation, nuclear emulsion stacks of several pellicles of NIFKI-BR2 type are used. The dimension of each Pellicle is  $18.7 \times 9.7 \times 0.06 \text{ cm}^3$ . The stacks were exposed to a beam of  $^{12}\text{C}$ -ions with momentum 4.5 A GeV/c at the Joint Institute for Nuclear Research (JINR) at Dubana, Russia. The search for inelastic interactions was carried out. A random sample of 485 events having  $N_h \geq 0$ , were picked up by using along the track doubly scanning method for detailed study. It may be pointed out that the interactions which were within 30  $\mu\text{m}$  from the top or bottom surface of the nuclear emulsion were not considered for the final investigation. Secondary tracks produced in each event were classified in to shower, grey, black, and heavily ionized tracks in a disintegration and are denoted by  $N_s, N_g, N_b$ , and  $N_h (= N_b + N_g)$ , respectively. The details of the scanning procedure, selection criteria used for

collecting the experimental data and the method of measuring angles etc. may be found elsewhere [22-24]. The secondary tracks emerging from each interaction are classified in accordance with their ionization or normalized grain density ( $g^*$ ), range ( $L$ ) and velocity ( $\beta$ ) as shower, grey and black particles. Shower particles are freshly created, singly charged relativistic particles with  $g^*$  less than 1.4. These particles have relativistic velocity  $\beta \geq 0.7$ . They are mostly fast pions with a small mixture of Kaons and released protons from the projectile which have undergone an interaction. For the case of proton, kinetic energy ( $E_p$ ) should be less than 400 MeV. The number of charged shower particles in an event is represented by  $N_s$ . Particles with range  $L > 3 \text{ mm}$  and  $1.4 < g^* < 6.0$  are defined as grey particles. They have  $\beta$  in the range of  $0.3 < \beta < 0.7$ . These are generally knocked out protons of targets with kinetic energy in between 30 – 400 MeV, and traces of deuterons, tritons and slow mesons. The number of charged grey particles in an event is represented by  $N_g$ . Particles having  $L < 3 \text{ mm}$  from interaction vertex from and  $g^* > 6.0$ . This corresponds to  $\beta < 0.3$  and protons of kinetic energy less than 30 MeV. Most of these are produced due to evaporation of residual target nucleus. The number of charged black particles in an event is represented by  $N_b$ . Black and grey particles are mainly product of breakup of target nucleus. The number of heavily ionizing charged particles ( $N_h$ ) are equal to the sum of black ( $N_b$ ) and grey ( $N_g$ ) fragments i.e.  $N_h (= N_b + N_g)$ .

## 3. Experimental Results and Discussion

### 3.1 The normalized moments of the multiplicity distributions

The normalized moments of the multiplicity distributions of relativistic charged particles produced in high energy hadron-nucleus collisions has been investigated by several workers in the energy range  $\sim (6-800) \text{ GeV}$ . However, a little attention has been paid to study the normalized moments of the multiplicity distributions of the relativistic charged particles produced in nucleus-nucleus interactions

[8-9]. Thus, it is considered worthwhile to investigate the normalized moments of the distributions and their dependence on the mass of the projectile and the struck nucleus [17-21]. The normalized moments of the multiplicity distributions of relativistic charged particles is defined as

$$C_k = \langle N_s^k \rangle / \langle N_s \rangle^k \quad (1)$$

Where, k is a constant and can have any value 2, 3, 4 etc.

For studying the dependence of  $C_k$  on the size of the target nucleus, the values of  $C_2, C_3, C_4$  are calculated for the different groups of emulsion nuclei in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus interactions. The value of  $C_k$  obtained in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus reactions are listed in Table1:

**Table 1:** value of  $C_k$  obtained in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus reactions

Types of Interactions	$C_2$	$C_3$	$C_4$	References
$^{12}\text{C}$ -CNO	$1.41 \pm 0.10$	$2.42 \pm 0.23$	$4.73 \pm 0.58$	Present work
$^{28}\text{Si}$ -CNO	$1.41 \pm 0.01$	$2.44 \pm 0.007$	$4.41 \pm 0.003$	4
$^{12}\text{C}$ -AgBr	$1.27 \pm 0.06$	$2.20 \pm 0.13$	$3.09 \pm 0.24$	Present work
$^{12}\text{C}$ -AgBr	$1.29 \pm 0.26$	$2.48 \pm 0.08$	-	3
$^{28}\text{Si}$ -AgBr	$1.31 \pm 0.007$	$4.56 \pm 0.004$	$3.55 \pm 0.35$	Present work
P-Em	$1.44 \pm 0.002$	$2.55 \pm 0.04$	-	1
$^{12}\text{C}$ -Em	$1.43 \pm 0.05$	$2.51 \pm 0.12$	$5.15 \pm 0.33$	Present work
$^{12}\text{C}$ -Em	$1.58 \pm 0.007$	$2.67 \pm 0.13$	-	3
$^{28}\text{Si}$ -Em	$1.53 \pm 0.01$	$3.06 \pm 0.013$	$7.19 \pm 0.01$	Present work

It is interesting to note in the table that the values of  $C_k$  in 4.5 A GeV/c proton-nucleus and nucleus-nucleus reactions are almost constant within the statistical limits. However, the

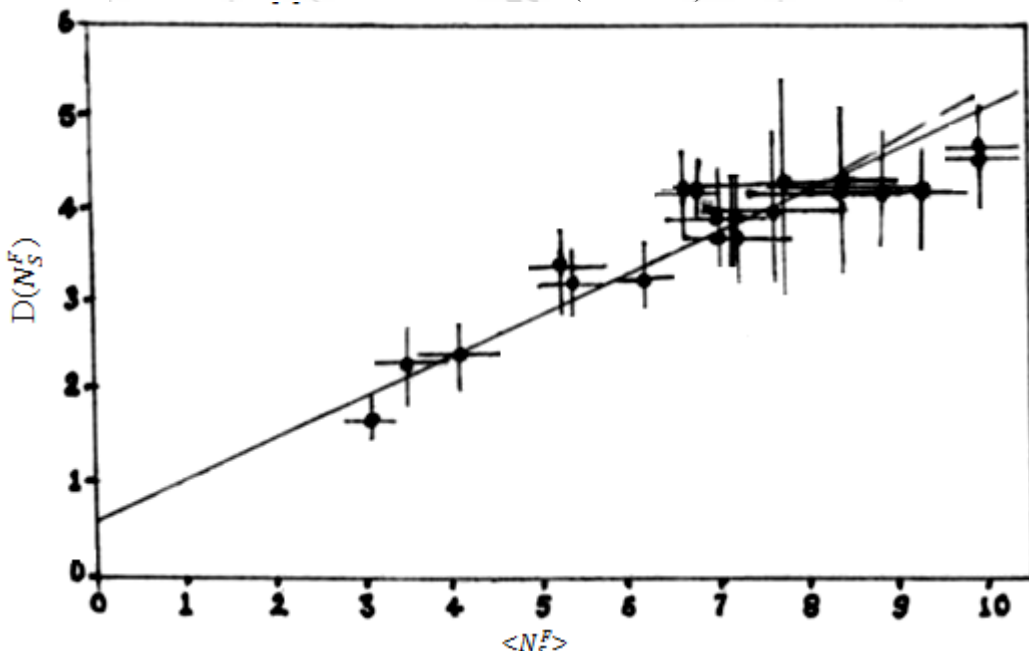
value of this parameter is observed to increase with the increasing values of k.

### 3.2 Characteristics of charged shower particles produced in the forward ( $0 < \theta < 90^\circ$ ) and backward ( $\theta > 90^\circ$ ) hemispheres:

Study of the characteristics of charged shower particles produced in the forward ( $0 < \theta < 90^\circ$ ) and backward ( $\theta > 90^\circ$ ) hemispheres, where  $\theta$  is the emission angle of relativistic charged particles in the laboratory system, has received considerable attention during the recent years in the nucleus-nucleus interactions [10-16]. It is reported [8] that the probability of the existence of leading particle is much higher in the forward hemisphere than in the backward hemisphere. Thus, the leading particle effect in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus reactions is investigated by studying the dependence of dispersion of relativistic charged particles in forward hemisphere,  $D(N_s^F) = (\langle N_s^F \rangle^2 - \langle N_s^F \rangle^2)^{1/2}$  on the average number of charged shower particles produced in forward hemisphere,  $\langle N_s^F \rangle$ . The variation of  $D(N_s^F)$  with  $\langle N_s^F \rangle$  is plotted in Fig. 1. It is evidently clear from the figure that the dispersion,  $D(N_s^F)$  grows linearly with  $\langle N_s^F \rangle$  in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus interactions. The experimental data obtained in the present work are found to satisfy the following relation.

$$D(N_s^F) = \alpha \langle N_s^F \rangle + \beta \quad (2)$$

The values of the coefficients  $\alpha$  and  $\beta$  are found to be  $(0.45 \pm 0.04)$  and  $(0.64 \pm 0.27)$  in 4.5 A GeV/c  $^{12}\text{C}$ -nucleus interactions. The parameter  $\alpha$  is called leading particle multiplicity. However, the value of  $\alpha$  is reported to be  $(1.04 \pm 0.13)$  in 4.5 A GeV/c  $^{28}\text{Si}$ -nucleus interactions [8].



**Figure 1:** Variation of  $D(N_s^F)$  with  $\langle N_s^F \rangle$  in  $^{12}\text{C}$

Thus, on comparing the value of  $\alpha$  obtained in the present work with those obtained in ref.9 and13. It may be concluded that the leading particle multiplicity strongly depends on the mass of the projectile.

On comparing these observations with those obtained in high energy hadron-nucleus [7] and proton-proton [6-9] collisions, it may be concluded that there is at least a

quantitative similarity in the mechanism of hadronization in the final stage of high energy hadron-nucleus as well as nucleus-nucleus reactions.

#### 4. Conclusions

- 1) The probability of the existence of leading particle is much higher in the forward hemisphere than in the backward hemisphere. This results indicates that the leading particle effect in 4.5 A GeV/c <sup>12</sup>C-nucleus reactions is investigated by studying the dependence of dispersion of relativistic charged particles in forward hemisphere.
- 2) The leading particle multiplicity strongly depends on the mass of the projectile.
- 3) In high energy hadron-nucleus and proton-proton collisions, there is at least a quantitative similarity in the mechanism of hadronization in the final stage of high energy hadron-nucleus as well as nucleus-nucleus reactions.

#### References

- [1] H.Khushnood et al: 6th Int. Conf. on Physics and Astrophysics of Quark Gluon Plasma, Dec.6-10,2010,Goa,India.
- [2] Praveen Prakash Shukla et al: Proc.DAE Symp.Nucl.Phys.Vol.59, 732 (2014).
- [3] Tauseef Ahmad et al: DAE Symp.of Nucl.Phys.,BARK,1991,India.
- [4] Dipak Ghosh et al: Czech.J.Phys.B 36,1358(1986).
- [5] Mahmoud Mohery: Cand. J. Phy 90 (12)1267, +1278, 2012.
- [6] B.P.Barik et al: Czech.J.Phys.B31,490(1981).
- [7] H.Khushnood et al: Can. J. Phys.68,67(1990).
- [8] M.Qasim Raza Khan: Ph.D. Thesis,AMU Aligarh(1988).
- [9] S.Sarfaz Ali: Ph.D. Thesis,Jamia Millia Islamia,New Delhi.
- [10] Tauseef Ahmad and M.Irfan: Phys.Rev.C46,1483(1992).
- [11] M. Saleem Khan et al: Nuovo Cim. A 109, 1623(1996).
- [12] A.Abdeslam: Phys.G.Nucl.Par.Phys.28, 1375(2002).
- [13] M.Man et al: Can.J.Phys. 62,230(1984).
- [14] L. Lohrman et al: Nuovo. Cim.,25,957(1962).
- [15] M. Saleem Khan et al: Il Nuovo Cim. A 108, 147(1995).
- [16] H.Khushnood et al: Can. J. Phys.61,1120(1983).
- [17] N. S. Chauhan et al: Indian Journal of Physics, 87 1263-1267(2013).
- [18] M. Mohery and M. Arafa : Internationa Journal of Modern Physics E, 20(2011) 1735 -1754.
- [19] T. Ahmad : IOSR Journal of Applied Physics (IOSR-JAP),Vol. 6, Issue 5 Ver. II, 15-19(2014)
- [20] M. Saleem Khan et al: J.Phys.Soc.Jpn.65, 801(1996).
- [21] M. Saleem Khan et al: Proc.DAE Symp.Nucl.Phys.Vol.56, 854(2011)
- [22] AALMT Collaboration: Yad. Faz. 22,736(1975).
- [23] AAGMT Collaboration: Nucl.Phys.B129, 205(1978).
- [24] I. Otterlund et al. Nucl.Phys. B142, 445 (1978).